

## 4. Übung QFT für Vielteilchensysteme

26.04.2012, 14:00-16:00, Seminarraum 138C

### 1. Getting familiar with the Density of States

2+3=5 points

Calculations of thermodynamic quantities, response functions and Feynman diagrams in QFT for condensed matter systems often require the evaluation of integrals or sums over all momenta  $\mathbf{k}$  (typically over the first Brillouin Zone). An important simplification of these  $\mathbf{k}$ -summations is possible, however, when the integrand  $\mathcal{F}$  is depending **on the energy**  $\varepsilon(\mathbf{k})$  only. In this case the integration/sum is best performed by using the energy  $\varepsilon$  as a variable. In the case of a cubic lattice of volume  $L^d$  in  $d$  dimensions we have for a given observable  $F$ :

$$F = \frac{1}{L^d} \sum_{\mathbf{k}} \mathcal{F}(\varepsilon_{\mathbf{k}}) = \frac{1}{(2\pi)^d} \frac{(2\pi)^d}{L^d} \sum_{\mathbf{k}} \mathcal{F}(\varepsilon_{\mathbf{k}}) \simeq \frac{1}{(2\pi)^d} \int d^d k \mathcal{F}(\varepsilon_{\mathbf{k}}) = \int d\varepsilon \mathcal{N}(\varepsilon) \mathcal{F}(\varepsilon) \quad (1)$$

where  $\mathcal{N}(\varepsilon)$ , i.e. the so-called **Density of States (DOS)**, which can be defined via comparison between the different equalities as

$$\mathcal{N}(\varepsilon) = \frac{1}{L^d} \sum_{\mathbf{k}} \delta(\varepsilon - \varepsilon_{\mathbf{k}}) \quad \text{or, for the continuous case,} \quad (2)$$

$$= \frac{1}{(2\pi)^d} \int d^d k \delta(\varepsilon - \varepsilon_{\mathbf{k}}). \quad (3)$$

- Calculate and plot the explicit expression for  $\mathcal{N}(\varepsilon)$  for non-interacting particles of mass  $m$  in the continuous case (i.e.,  $\varepsilon_{\mathbf{k}} = \frac{\hbar^2 k^2}{2m}$ ) in one, two and three dimensions. How do the corresponding Fermi surfaces look like in these cases?
- Consider the energy dispersion  $\varepsilon(\mathbf{k}) = -2t \sum_{i=1}^d \cos(k_i a)$  derived in the last exercise for the  $d$ -dimensional Hubbard-model. Try to plot (numerically) the DOS  $\mathcal{N}(\varepsilon)$  for the cases  $d = 1, 2, 3$ . Which are the most prominent features of these DOS functions and at which energies  $\varepsilon$  they occur? How would the corresponding Fermi surfaces look like for the case  $d = 1, 2$ , e.g. if one has an average density of one electron per site (*half-filled system*)? For further information see also: A. Georges, G. Kotliar, et.al., Rev.Mod.Phys. 68 (Appendix A).

### 2. Calculations of the Lindhard function (I)

2+3=5 points

In perturbation theory the state of a system, at first order in the perturbation, is given by:

$$|\psi_{\mathbf{k}}\rangle = |\psi_{\mathbf{k}}^0\rangle + \sum_{\mathbf{k}'} \frac{|\psi_{\mathbf{k}'}^0\rangle \langle \psi_{\mathbf{k}'}^0 | V | \psi_{\mathbf{k}}^0 \rangle}{\varepsilon_{\mathbf{k}} - \varepsilon_{\mathbf{k}'}}. \quad (4)$$

- Suppose that the perturbation  $V$  is generated by an external<sup>1</sup> charge distribution:

$$\langle \mathbf{r} | V | \mathbf{r}' \rangle = -e \delta(\mathbf{r} - \mathbf{r}') \phi(\mathbf{r}), \quad (5)$$

and, assuming that the wave functions  $\psi(\mathbf{r}) = \langle \mathbf{r} | \psi_{\mathbf{k}}^0 \rangle$  of the unperturbed system are plain waves, express the matrix element  $\langle \psi_{\mathbf{k}}^0 | V | \psi_{\mathbf{k}+\mathbf{q}}^0 \rangle$  in terms of the Fourier transform  $\phi(\mathbf{q})$  of the electrostatic potential.

b) Express the distribution of the electronic charge density as:

$$\rho(\mathbf{r}) = -e \sum_{\mathbf{k}} f_{\mathbf{k}} |\psi_{\mathbf{k}}(\mathbf{r})|^2 = \rho^0 + \rho^{\text{ind}}(\mathbf{r}) \quad (6)$$

(where  $f_{\mathbf{k}}$  is the equilibrium Fermi distribution),  $\rho^0$  is the density of the unperturbed system and  $\rho^{\text{ind}}(\mathbf{r})$  is the variation of the density at *first order* in the perturbation. Show that the Fourier transform of the charge induced to first order in the potential  $\phi$  is given by

$$\rho^{\text{ind}}(\mathbf{q}) = -e^2 \int \frac{d^3k}{4\pi^3} \frac{f_{\mathbf{k}-\frac{1}{2}\mathbf{q}} - f_{\mathbf{k}+\frac{1}{2}\mathbf{q}}}{\hbar^2(\mathbf{k} \cdot \mathbf{q}/m)} \phi(\mathbf{q}). \quad (7)$$

In which limit does the Lindhard screening approach the Thomas-Fermi one?

<sup>1</sup> The perturbation is generated by an *external* distribution of charge, therefore it can be treated here as a *one particle* potential.

### 3. Screened and unscreened Coulomb Potentials

0 points

(Additional exercise for the Thomas-Fermi model which will not be graded, but just discussed by the tutor.)

a) From the integral representation of the delta function,

$$\delta(\mathbf{r}) = \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{r}} \quad (8)$$

and the fact that the Coulomb potential  $\phi(\mathbf{r}) = -e/r$  satisfies Poisson's equation,

$$-\nabla^2 \phi(\mathbf{r}) = -4\pi e \delta(\mathbf{r}), \quad (9)$$

show that the electronic pair potential,  $V(\mathbf{r}) = -e\phi(\mathbf{r}) = e^2/r$ , can be written in the form

$$V(\mathbf{r}) = \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{r}} V(\mathbf{k}), \quad (10)$$

where the Fourier transform  $V(\mathbf{k})$  is given by

$$V(\mathbf{k}) = \frac{4\pi e^2}{k^2} \quad (11)$$

b) Show that the Fourier transf. of the screened Coulomb interaction  $V_s(\mathbf{r}) = (e^2/r)e^{-k_{TF}r}$  is

$$V_s(\mathbf{k}) = \frac{4\pi e^2}{k^2 + k_{TF}^2} \quad (12)$$

by substituting (12) into the Fourier integral

$$V_s(\mathbf{r}) = \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{r}} V_s(\mathbf{k}), \quad (13)$$

and evaluating that integral in spherical coordinates (*Hint*: The radial integral is best done as a contour integral.). Finally, deduce from (12) that  $V_s(\mathbf{r})$  satisfies

$$(-\nabla^2 + k_{TF}^2) V_s(\mathbf{r}) = 4\pi e^2 \delta(\mathbf{r}) \quad (14)$$