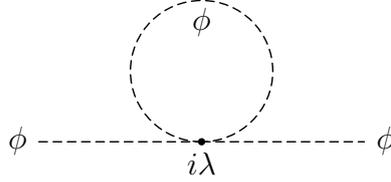


43. Working in dimensional regularisation as regularisation scheme in D -dimensions, in every closed loop with a four-particle vertex there occurs the so-called one-point function A_0 , with e.g. the matrix element $\mathcal{M} = i\lambda A_0(m)$, depicted by the Feynman graph



It is defined as

$$A_0(m) = \frac{(2\pi\mu)^{4-D}}{i\pi^2} \int d^D k \frac{1}{k^2 - m^2 + i\epsilon}. \quad (57)$$

Using the Wick rotation $k^0 \rightarrow ik_E^0$ the integration path in the complex k^0 plane will be rotated by $\pi/2$ in order to get Euklidian coordinates. Use the Bogolubov-Schwinger parametrisation

$$\frac{1}{A} = \int_0^\infty dx e^{-xA} \quad A > 0. \quad (58)$$

Show that the intermediate result has the form

$$\frac{i}{16\pi^2} A_0(m) = -i \frac{\mu^{4-D}}{(2\pi)^D} \int_0^\infty d\alpha e^{-\alpha m^2} \int d^D k_E e^{-\alpha k_E^2}. \quad (59)$$

For the integration over the D -dimensional momentum space work with spherical coordinates. For that step these formulas are helpful:

$$\int d\Omega_D = \frac{2\pi^{D/2}}{\Gamma(D/2)} \quad \text{and} \quad \int_0^\infty dx x^a e^{-bx^2} = \frac{\Gamma\left(\frac{a+1}{2}\right)}{2b^{\frac{a+1}{2}}} \quad b > 0, \quad a \in N_u. \quad (60)$$

The remaining parameter integral is partially integrated (the exponent of α becomes $1 - D/2$). By using

$$\int_0^\infty dx x^a e^{-bx} = \Gamma(a+1) b^{-a-1} \quad (61)$$

and also $\epsilon = 4 - D$ verify the result

$$\frac{i}{16\pi^2} A_0(m) = \frac{2im^2}{(4\pi)^2} \frac{\Gamma\left(\frac{\epsilon}{2}\right)}{2 - \epsilon} e^{\frac{\epsilon}{2} \left(\ln 4\pi + \ln \frac{\mu^2}{m^2} \right)}. \quad (62)$$

Perform an expansion in ϵ , $\left(\Gamma\left(\frac{\epsilon}{2}\right) = \frac{2}{\epsilon} - \gamma + \mathcal{O}(\epsilon)\right)^1$ and show that the final result is

$$A_0(m) = m^2 \left(\frac{2}{\epsilon} - \gamma + 1 + \ln 4\pi + \ln \frac{\mu^2}{m^2} \right) + \mathcal{O}(\epsilon). \quad (63)$$

¹ $\gamma = 0.5772157\dots$ Euler-Mascheroni constant