

Black Holes II — Exercise sheet 3

(13.1) Schwarzschild black holes do not bifurcate

We saw last semester that a black hole binary system can lead to a single black hole final state, after emitting an appreciable amount of gravitational radiation in the inspiralling, merger and ringdown phases. The endstate was a Kerr black hole with mass M and angular momentum J . We also saw that a Kerr black hole is unstable against superradiance, which lowers its angular momentum and mass, while maintaining (or increasing) its horizon area. The endstate of this chain of instabilities is a Schwarzschild black hole. Prove that this endstate is classically stable against bifurcation. In other words, show that a Schwarzschild black hole cannot decay into two (Schwarzschild or Kerr) black holes.¹

(13.2) Raychaudhuri equation for geodesic null congruences

Consider a congruence of null geodesics with tangent vector field k^a and deviation vector η^a (in 4 spacetime dimensions). Derive the analog of the Raychaudhuri equation for the expansion θ . You should find

$$k^a \nabla_a \theta = -\frac{1}{2} \theta^2 - \sigma_{ab} \sigma^{ab} + \omega_{ab} \omega^{ab} - R_{ab} k^a k^b$$

Explain why there is a factor $\frac{1}{2}$ instead of $\frac{1}{3}$.

(13.3) Prove Your Own Singularity Lemma

Given a geodesic null congruence with vanishing twist ($\omega_{ab} = 0$) and assuming the null energy condition (so that $R_{ab} k^a k^b \geq 0$) show the following lemma: If the expansion is negative, $\theta = \theta_0 < 0$, at some point on a geodesic in the congruence then $\theta \rightarrow -\infty$ along that geodesic within the affine distance $\lambda \leq 2/|\theta_0|$.

These exercises are due on March 25th 2014.

¹Unless, of course, something provides sufficient energy to trigger such a process. But when we use the word “decay” we usually mean “decay all by itself, without some external trigger”. Sidenote: the black holes that we observe have a surrounding accretion disk, which actually tends to bring the rotating black hole towards extremality; the processes described in this exercise work in the other direction and apply to black holes in vacuum.

Hints:

- Use Hawking's area theorem. Use energy conservation. Perhaps have a look at exercises (10.2) and (10.3). This is a short exercise!
- Note that the orthogonality condition $k^a \eta_a = 0$ is not sufficient to determine uniquely η^a since the displacement vector may now have a component parallel to k^a . Thus, the displacement vector η^a orthogonal to the tangent vector k^a specifies only a co-dimension 2 parameter family of geodesics (in our 4 spacetime dimensions this is a 2-parameter family). Therefore, you must specify one further condition to fix the displacement vector η^a . It is convenient to introduce another null-vector l^a with the properties $l^a l_a = 0$, $l^a k_a = -1$ and $k^a \nabla_a l^b = 0$ (so if k^a is an ingoing null vector l^a is an outgoing one, with some convenient normalization). Then you may impose the condition

$$\eta^a l_a = 0$$

Instead of the "spatial metric" $h_{ab} = g_{ab} + t_a t_b$ introduce now the projector

$$P_{ab} = g_{ab} + k_a l_b + l_a k_b$$

that projects onto the required co-dimension 2 subspace of the tangent space, $P^a_b \eta^b = \eta^a$. The tensor field $B_{ab} = \nabla_b k_a$ must also be projected,

$$\hat{B}^a_b := P^a_c B^c_d P^d_b$$

After showing that $k^a \nabla_a \eta^b = \hat{B}^b_a \eta^a$ (remember that $k^a \nabla_a \eta^b = \eta^a \nabla_a k^b$) make the decomposition

$$\hat{B}^a_b = \frac{1}{2} \theta P^a_b + \sigma^a_b + \omega^a_b$$

and think about the correctness of the factor $\frac{1}{2}$. The rest is completely analog to the derivation of the Raychaudhuri equations for timelike geodesic congruences. This is a long exercise!

- This lemma is the light-like analog of the lemma we proved for time-like geodesic congruences during the lectures. You need the result for the Raychaudhuri equation for geodesic null congruences spelled out in exercise (13.2), and otherwise proceed analog to the lectures.