

# 1. Problem Set, 12.10.2020, 1h<sup>00</sup>

## 1.1. IF PHOTON HAS MASS

$$\vec{E} = -\vec{\nabla}\Phi, \quad \Delta\Phi = -\frac{\rho}{\epsilon_0} + \frac{\rho}{L^2}$$

in Space between the shells  $\rho = 0$

$$\rightarrow \Delta\Phi = \frac{\rho}{L^2}, \quad \psi(r) = \frac{u(r)}{r} \rightarrow \text{ANSATZ}$$

$$\Delta\Phi = \frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{d}{dr} \psi(r) \right) = \frac{u(r)}{L^2} \quad 0 \leq r_1 < r < r_2$$

Spherical symmetry

$$\frac{1}{r^2} \frac{d}{dr} \left[ r^2 \left( \frac{d}{dr} u(r) \frac{1}{r} - u(r) \frac{1}{r^2} \right) \right] = \frac{u(r)}{L^2 r}$$

$$\frac{1}{r^2} \frac{d}{dr} \left[ r \frac{d}{dr} u(r) - u(r) \right] = \frac{u(r)}{L^2 r}$$

$$\frac{1}{r} \left[ \frac{d}{dr} u(r) + r \frac{d^2}{dr^2} u(r) - u(r) \right] = \frac{u(r)}{L^2}$$

$$\frac{d^2}{dr^2} u(r) = \frac{u(r)}{L^2}$$

$$u(r) = e^{\alpha r} \rightarrow \alpha^2 u(r) = \frac{u(r)}{L^2}$$

$$\alpha^2 = \frac{1}{L^2} \rightarrow \alpha = \pm \frac{1}{L}$$

$$\psi(r) = \frac{1}{r} \left[ A e^{r/L} + B e^{-r/L} \right]$$

$$V = \frac{1}{r_1} \left[ A e^{r_1/L} + B e^{-r_1/L} \right] \rightarrow r_1 V = A e^{r_1/L} + B e^{-r_1/L}$$

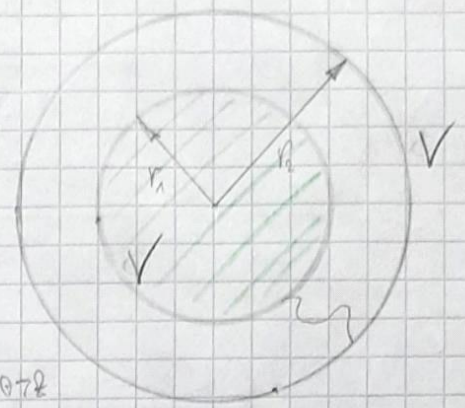
$$B e^{-r_1/L} = r_1 V - A e^{r_1/L} \rightarrow B = r_1 V e^{r_1/L} - A e^{2r_1/L}$$

$$r_2 V = A e^{r_2/L} + B e^{-r_2/L} = A e^{r_2/L} + e^{-r_2/L} \left( r_1 V e^{r_1/L} - A e^{2r_1/L} \right)$$

$$r_2 V - r_1 V e^{r_1/L} e^{-r_2/L} = A \left( e^{r_2/L} - e^{-r_2/L} e^{2r_1/L} \right)$$

$$A = \frac{r_2 V - r_1 V e^{r_1/L} e^{-r_2/L}}{e^{r_2/L} - e^{-r_2/L} e^{2r_1/L}} \rightarrow B = r_1 V e^{r_1/L} - e^{2r_1/L} \cdot$$

$$\cdot \frac{r_2 V - r_1 V e^{r_1/L} e^{-r_2/L}}{e^{r_2/L} - e^{-r_2/L} e^{2r_1/L}}$$



$$\begin{aligned} \phi(r) &= \frac{V}{r} \left[ e^{\frac{r_1}{L}} \frac{r_2 \sqrt{L} - r_1 \sqrt{L} e^{\frac{r_1}{L}} e^{-\frac{r_2}{L}}}{e^{\frac{r_1}{L}} - e^{-\frac{r_2}{L}} e^{\frac{2r_1}{L}}} + e^{-\frac{r_1}{L}} \left( r_1 \sqrt{L} e^{\frac{r_1}{L}} - e^{\frac{2r_1}{L}} \right) \right. \\ &\quad \left. + \frac{r_2 \sqrt{L} - r_1 \sqrt{L} e^{\frac{r_1}{L}} e^{-\frac{r_2}{L}}}{e^{\frac{r_1}{L}} - e^{-\frac{r_2}{L}} e^{\frac{2r_1}{L}}} \right] \\ &= \frac{V}{r} \left[ \frac{r_2 e^{\frac{r_1}{L}} - r_1 e^{\frac{r_1}{L}} e^{-\frac{r_2}{L}} e^{\frac{r_1}{L}}}{\text{Nenner}} + \frac{e^{-\frac{r_1}{L}} r_1 e^{\frac{r_1}{L}} \left( e^{\frac{r_1}{L}} - e^{-\frac{r_2}{L}} e^{\frac{2r_1}{L}} \right) - e^{-\frac{r_1}{L}} e^{\frac{2r_1}{L}}}{\text{Nenner}} \right. \\ &\quad \left. + (r_2 - r_1) e^{\frac{r_1}{L}} e^{-\frac{r_2}{L}} \right] \end{aligned}$$

$$= \frac{V}{r} \left[ r_2 e^{\frac{r_1}{L}} - r_1 e^{\frac{r_1 - r_2 + r}{L}} + r_1 e^{\frac{r_1 + r_2 - r}{L}} - r_1 e^{\frac{3r_1 - r_2 - r}{L}} - r_2 e^{\frac{2r_1 - r}{L}} + r_1 e^{\frac{3r_1 - r_2 - r}{L}} \right] \frac{1}{\text{Nenner}}$$

$$= \frac{V}{r} \left[ r_2 \frac{e^{\frac{r_1}{L}} - e^{\frac{2r_1}{L}} e^{-\frac{r_2}{L}}}{e^{\frac{r_1}{L}} - e^{-\frac{r_2}{L}} e^{\frac{2r_1}{L}}} + r_1 \frac{e^{\frac{r_1 + r_2 - r}{L}} - e^{\frac{r_1 - r_2 + r}{L}}}{e^{\frac{r_1}{L}} - e^{-\frac{r_2}{L}} e^{\frac{2r_1}{L}}} \right]$$

$$= \frac{V}{r} \left[ \frac{r_2}{r} \frac{e^{\frac{r-r_1}{L}} - e^{\frac{r_1+r}{L}}}{e^{\frac{r_2-r_1}{L}} - e^{\frac{r_1-r_2}{L}}} + \frac{r_1}{r} \frac{e^{\frac{r_2-r}{L}} - e^{\frac{r-r_2}{L}}}{e^{\frac{r_2-r_1}{L}} - e^{\frac{r_1-r_2}{L}}} \right]$$

$$= \frac{V}{r} \left[ \frac{r_2}{r} \frac{\sinh\left(\frac{r-r_1}{L}\right)}{\sinh\left(\frac{r_2-r_1}{L}\right)} + \frac{r_1}{r} \frac{\sinh\left(\frac{r_2-r}{L}\right)}{\sinh\left(\frac{r_2-r_1}{L}\right)} \right]$$

$$\alpha := \frac{r_2 - r_1}{L} \rightarrow \phi(r) = \frac{V}{r} \left[ \frac{r_2}{r} \frac{\sinh\left(\frac{r-r_1}{L}\right)}{\sinh \alpha} + \frac{r_1}{r} \frac{\sinh\left(\frac{r_2-r}{L}\right)}{\sinh \alpha} \right]$$

$$\begin{aligned} \vec{E} &= -\vec{\nabla} \phi(r) = \hat{r} \nabla \phi(r) = \\ &= \frac{V \hat{r}}{\sinh \alpha} \left[ r_2 \left[ \frac{\sinh\left(\frac{r-r_1}{L}\right)}{r^2} - \frac{\cosh\left(\frac{r-r_1}{L}\right)}{rL} \right] + r_1 \left[ \frac{\sinh\left(\frac{r_2-r}{L}\right)}{r^2} + \frac{\cosh\left(\frac{r_2-r}{L}\right)}{rL} \right] \right] \end{aligned}$$

$$\begin{aligned} \text{b) } -\nabla^2 \phi &= \vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} = \frac{Q_{\text{end}}}{L^2} \\ \int_V \vec{\nabla} \cdot \vec{E} dV &= \int_V d\vec{S} \cdot \vec{E} = \int_V \frac{\rho}{\epsilon_0} dV = \frac{Q_{\text{end}}}{\epsilon_0} = \frac{1}{L^2} \int_V \phi(r) dV \end{aligned}$$

$$\int_{\partial V} d\vec{s} \cdot \vec{E} = \epsilon_0 \int_V \nabla \cdot \vec{E} = \frac{Q_1}{\epsilon_0} = \frac{1}{L^2} \int dV \rho(r)$$

$$\frac{1}{L^2} \frac{V}{\sinh \alpha} \left\{ r_2 \left[ -\frac{1}{r_1 L} \right] + \frac{1}{r_1} \left[ \frac{\sinh \left( \frac{L-r_1}{L} \right)}{r_1} + \frac{\cosh \left( \frac{r_2-r_1}{L} \right)}{r_1/L} \right] \right\}$$

$$= \frac{Q_1}{\epsilon_0} = \frac{1}{L^2} V \frac{4\pi r_1^3}{3}$$

$$\frac{4\pi r_1^3 V}{3 \sinh \alpha} \left\{ -\frac{r_2}{r_1 L} + \frac{\sinh \alpha}{r_1} + \frac{\cosh \alpha}{L} \right\} = \frac{Q_1}{\epsilon_0} = \frac{V 4\pi r_1^3}{3 L^2}$$

$$4\pi r_1^3 V \left\{ \frac{1}{r_1} - \frac{r_2 \operatorname{csch} \alpha}{r_1 L} + \frac{\cosh \alpha}{L} \right\} = \frac{Q_1}{\epsilon_0} = \frac{V 4\pi r_1^3}{3 L^2}$$

$$Q_1 = 4\pi \epsilon_0 r_1^2 V \left\{ \frac{1}{r_1} + \frac{\cosh \alpha}{L} - \frac{r_2}{r_1 L} \operatorname{csch} \alpha \right\} + \frac{4\pi \epsilon_0}{3 L^2} r_1^3 V$$

a)  $L \rightarrow \infty \Rightarrow \alpha \rightarrow 0$

$\operatorname{csch} x \sim \frac{1}{x} - \frac{x}{6}$ ,  $\cosh x \sim \frac{1}{x} + \frac{x}{3}$

$$\Rightarrow Q_1 \sim 4\pi \epsilon_0 r_1^2 V \frac{(r_2 - r_1)(2r_1 + r_2)}{6L^2 r_1} + \frac{4\pi \epsilon_0}{3 L^2} r_1^3 V$$

$$= \frac{2\pi \epsilon_0}{3} \frac{r_1 V}{L^2} \left( \frac{r_2}{r_1} \right)^2 \left( 1 + \frac{r_1}{r_2} \right)$$

## 1.2 CONSEQUENCES OF GAUGE CHOICES

a)  $\nabla \cdot \vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int d^3 r' \vec{j}(\vec{r}') \cdot \nabla \frac{1}{|\vec{r} - \vec{r}'|} = \frac{1}{3}$

$\nabla \frac{1}{r} = -\frac{\vec{r}}{r^3} \Rightarrow \nabla \frac{1}{|\vec{r} - \vec{r}'|} = -\frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3}$ ,  $\nabla' \frac{1}{|\vec{r} - \vec{r}'|} = +\frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3}$

$\frac{1}{3} = -\frac{\mu_0}{4\pi} \int d^3 r' \vec{j}(\vec{r}') \cdot \nabla' \frac{1}{|\vec{r} - \vec{r}'|} = 0$  for steady currents

$$= -\frac{\mu_0}{4\pi} \int d^3 r' \nabla' \left( \frac{\vec{j}(\vec{r}')}{|\vec{r} - \vec{r}'|} \right) = \int d^3 r' \frac{\nabla' \cdot \vec{j}(\vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

Kontinuitätsgl.:  $\frac{\partial \rho}{\partial t} + \nabla \cdot \vec{j} = 0$

•  $-\frac{\mu_0}{4\pi} \int d\vec{S} \cdot \vec{j}(\vec{r}') \frac{1}{|\vec{r}-\vec{r}'|} \rightarrow$  surface integral vanishes at infinity, where we presume  $\vec{j}(\vec{r}) \rightarrow 0$

b) Biot-Savart:  $\nabla \cdot \vec{B} = 0$  &  $\nabla \times \vec{B} = \mu_0 \vec{j}$

$\rightarrow \vec{B}(\vec{r}) = \nabla \times \frac{\mu_0}{4\pi} \int d^3r' \frac{\vec{j}(\vec{r}')}{|\vec{r}-\vec{r}'|}$

(Helmholtz theorem)

Biot-Savart formula in conventional form

$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int d^3r' \frac{\vec{j}(\vec{r}') \times (\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3}$  valid when

$\nabla \cdot \vec{B} = 0$  &  $\nabla \times \vec{B} = \mu_0 \vec{j}$   $\rightarrow$  necessary for HT

Compare with proposed formula

$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int d^3r' \frac{\vec{j}(\vec{r}') \times (\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3}$

$\rightarrow \nabla \cdot \vec{A} = 0$  &  $\nabla \times \vec{A} = \vec{B}$

$\rightarrow$  COULOMB GAUGE

# 1.3 AHARONOV-BOHM EFFECT

a) Coulomb gauge:  $\vec{\nabla} \cdot \vec{A} = 0$  &  $\vec{\nabla} \times \vec{A} = \vec{B} = B \vec{e}_z$   
 Due to sym.  $\vec{A} = A \vec{e}_\phi$  &  $A = A(r)$

magnetic flux:  $\Phi_0 = \int_S \vec{dS} \cdot \vec{B} = \int_S \vec{dS} \cdot \vec{\nabla} \times \vec{A} = \int_{\partial S} d\vec{\ell} \cdot \vec{A}$

$(r > R)$   $\Phi_0 = B R^2 = 2\pi r A(r)$

$(r < R)$   $\Phi_0 = B r^2 = 2\pi r A(r)$

$$A(r) = \begin{cases} Br/2, & r < R \\ \frac{BR^2}{2r}, & r \geq R \end{cases} \longrightarrow A = A(r) \vec{e}_\phi$$

$r = R$   $\begin{cases} Br/2 \\ BR^2/2R = \frac{BR}{2} \end{cases}$  ✓

$\vec{\nabla} \cdot \vec{A} = 0 \longrightarrow$  Coulomb gauge

b)  $\vec{A}' = \vec{A} + \vec{\nabla} \chi = \vec{A} + \vec{\nabla} \left( -\frac{\Phi_0 \phi}{2\pi} \right)$ ,  $\Phi_0 = \text{const for } r > R$

$\vec{\nabla} = \vec{e}_r \frac{\partial}{\partial r} + \frac{1}{r} \vec{e}_\phi \frac{\partial}{\partial \phi} + \vec{e}_z \frac{\partial}{\partial z}$

$$\vec{A}' = \vec{A} - \frac{1}{2\pi r} \Phi_0 \vec{e}_\phi = \vec{A} - \frac{BR^2}{2\pi r} \vec{e}_\phi$$

$$\vec{A}' = \begin{cases} \left( Br/2 - \frac{BR^2}{2r} \right) \vec{e}_\phi, & r < R \\ \left( \frac{BR^2}{2r} - \frac{BR^2}{2r} \right) \vec{e}_\phi = 0, & r \geq R \end{cases}$$

c)  $\vec{B}' = \vec{\nabla} \times \vec{A}' = \vec{\nabla} \times \left( \vec{A} + \vec{\nabla} \chi \right)$  (for  $r < R$ )

$$\vec{B}' = \frac{\Phi_0}{2\pi} \vec{\nabla} \times \left( \frac{\vec{e}_\phi}{r} \right)$$

for  $r \neq 0$ :  $\vec{\nabla} \times \left( \frac{1}{r} \vec{e}_\phi \right) = \frac{1}{r} \frac{\partial}{\partial r} \left( r \cdot \frac{1}{r} \right) \vec{e}_z = 0$

Case  $r=0$   $\int_S \vec{dS} \cdot \left[ \vec{\nabla} \times \left( \frac{\vec{e}_\phi}{r} \right) \right] = \oint_{\partial S} d\vec{\ell} \cdot \frac{\vec{e}_\phi}{r} = \int_0^{2\pi} d\phi = 2\pi$

$$\Rightarrow \vec{\nabla} \times \left( \frac{\vec{e}_z}{r} \right) = \frac{J(r)}{r} \vec{e}_z$$

$$\Rightarrow \underline{\vec{B}'} = \vec{B} - \frac{\Phi_0}{2\pi} \frac{J(r)}{r} \vec{e}_z$$