

## 2. Übung QFT für Vielteilchen-Systeme

*22.03.2012, 14:00-16:00, Seminarraum 138C*

### 1. Bose-distribution function

*1+3 Punkte*

Consider a system of  $N$  non-interacting harmonic oscillators with frequencies  $\omega_i$ , i.e.,

$$\hat{\mathcal{H}} = \sum_{i=1}^N \hat{\mathcal{H}}_i = \sum_{i=1}^N \hbar\omega_i \left(\hat{n}_i + \frac{1}{2}\right) \quad \hat{n}_i = \hat{a}_i^\dagger \hat{a}_i. \quad (1)$$

a)  $Z$ , the so called partition function (“Zustandssumme”), is defined as  $Z = \text{Tr} \left( e^{-\beta \hat{\mathcal{H}}} \right)$ . Show that the total energy of a system, given as  $\langle \hat{\mathcal{H}} \rangle$  (see exercise 1), can be calculated as

$$E = \langle \hat{\mathcal{H}} \rangle = -\frac{\partial \ln Z}{\partial \beta}. \quad (2)$$

b) Calculate the partition function  $Z$  for the Hamiltonian of Eq. (1) and show that the total energy can be written as

$$E = E_0 + \sum_{i=1}^N \hbar\omega_i b(\omega_i), \quad (3)$$

where  $b$  is the Bose-distribution function

$$b(\omega_i) = \frac{1}{e^{\beta \hbar \omega_i} - 1}, \quad (4)$$

and  $E_0$  is the zero-point energy of  $N$  harmonic oscillators.

Hints: Since the oscillators are non-interacting the eigenstates of the system (which one uses for calculating the trace), can be written as the product of one-oscillator eigenstates, i.e.,  $|\{n_i\}\rangle = |n_1\rangle \otimes \dots \otimes |n_N\rangle$ , where  $n_i = 0 \dots \infty$  is the occupation of the  $i$ -th oscillator.

### 2. Fermi-distribution function

*3+3\* = 6 Punkte*

The entropy  $S$  of a system with total energy  $E$  and  $N$  particles is given by

$$S(E, N) = -\ln W(E, N), \quad (5)$$

where (for a discrete system)  $W(E, N)$  is the number of configurations for given values of  $E$  and  $N$ .

The (discrete) one-particle energy-levels  $\varepsilon_i$  (non interacting) of the system are  $g_i$ -fold degenerate ( $g_i \gg 1$ ). For large systems ( $N \gg 1$ )  $W$  can be approximated by

$$W = \prod_i w_i, \quad (6)$$

where  $w_i$  is the number of possible configurations of  $\bar{n}_i$  identical fermions in  $g_i$  degenerate states. Here,  $\bar{n}_i$  is the (still unknown!) most probable occupation of the ( $g_i$ -fold degenerate) energy level

$\varepsilon_i$  and  $g_i > \bar{n}_i$ . The  $\bar{n}_i$  can be computed by finding the extremum of the entropy with respect to  $\bar{n}_i$  considering the constraints

$$E = \sum_i \varepsilon_i \bar{n}_i \quad (7)$$

$$N = \sum_i \bar{n}_i. \quad (8)$$

- a) Calculate  $w_i$  for a given (integer) number  $\bar{n}_i$  of fermions and degeneracy  $g_i (> \bar{n}_i)$  of the energy level  $\varepsilon_i$ . Hint: Start considering  $g_i = 2, 3, 4, \dots$ . How many possible configurations exist for  $\bar{n}_i = 1, 2, 3, \dots$  independent fermions? Finally, calculate the number of configurations for a general  $g_i$  and  $\bar{n}_i$ .
- b) Calculate the extremum of the entropy  $S = -\ln W$ , with  $W$  as given in Eq. (6) with respect to  $\bar{n}_i$ . In order to include the constraints (7) and (8) use the Lagrange multipliers  $\beta$  for the energy and  $\alpha$  for the particle number.

Hints: Make use of the Stirling formula for the factorial of a large number  $N$ :

$$\ln N! \sim N(\ln N - 1). \quad (9)$$

For further information see also Ref. 1.

[1] Kerson Huang, “Statistical Mechanics”, chapter 4.3 and chapter 7.

\* Bonuspoints