
6. Exercise on QFT for many-body systems

Sommersemester 2018

TUTORIUM: Friday, 22.06.2018.

12. Van Hove singularities

2+2+1+2 = 5+2* points*

Consider the dispersion relation (single-particle energy states) for electrons on a simple hypercubic lattice in d dimensions, with only nearest-neighbor hopping:

$$\varepsilon_{\mathbf{k}} = -2t \sum_{i=1}^d \cos k_i, \quad (1)$$

with the hopping amplitude t and the lattice constant $a = 1$. The density of single-particle states in this system is then given by

$$N(\epsilon) = \frac{1}{(2\pi)^d} \int_{-\pi}^{\pi} d^d k \delta(\epsilon - \varepsilon_{\mathbf{k}}). \quad (2)$$

In the first exercise you have calculated numerically and then plotted these densities of states for $d = 1, 2, 3$. Here, the singular structures (divergences, cusps) of these functions should be analyzed analytically.

- a) Calculate $N(\epsilon)$ for $d = 1$ explicitly and determine the interval $[\epsilon_1, \epsilon_2]$ on which $N(\epsilon) \neq 0$. Moreover, identify the values ϵ^* where D diverges, i.e. where $N(\epsilon^*) = \infty$. From which points \mathbf{k}^* in the dispersion relation originate these divergences? Show that the divergences can be reproduced by taking into account only the contributions from these \mathbf{k}^* -points. (*Hint: Replace $\varepsilon_{\mathbf{k}}$ in Eq. (2) by a corresponding Taylor-expansion around these points up to second order.*)
- b) For $d = 2$ one can show that $N(\epsilon)$ is essentially given by a complete elliptic integral of the first kind. Here, however, only the singular contributions to $N(\epsilon)$ should be analyzed. As in the one-dimensional case a singular contribution originates from stationary points in the dispersion relation. Determine the kind of stationary point (i.e., maximum, minimum or saddle point) which generates this so-called Van Hove singularity in the two-dimensional DOS and determine the singular contribution to $N(\epsilon)$ by expanding $\varepsilon_{\mathbf{k}}$ around corresponding stationary point in Eq. (2) as for the one-dimensional case in a).
- c) Try to predict how the singular behavior of the DOS evolves with the dimensions of the system for $d \geq 3$.
- d) Finally, consider the limit $d \rightarrow \infty$. In this case, one has to rescale the hopping amplitude as $t \rightarrow \frac{t}{\sqrt{d}}$, in order to render the total energy of the system as well as the second moment (standard deviation) of the density of state finite. Show that $N_{\infty}(\epsilon)$ is proportional to a Gaußdistribution.

13. Magnetic susceptibilities in d dimensions

1.5+1.5+1+1=5 points

Consider a system of non-interacting electrons on a (hyper)cubic lattice whose energy dispersion is given by Eq. (1).

- a) Compute the *magnetic* susceptibility, i.e. the Fourier transform of the spin-spin response function $\langle T_\tau S_z(\mathbf{r}_i, \tau) S_z(0, 0) \rangle$, for the frequency $\Omega_m = 0$ (static susceptibility), and for the two momenta $\mathbf{Q} = (0, 0, 0, \dots)$ (ferromagnetic susceptibility) and $\mathbf{Q} = (\pi, \pi, \pi, \dots)$ (antiferromagnetic susceptibility).
- b) Determine the leading divergences of the ferromagnetic and the antiferromagnetic susceptibilities for $T \rightarrow 0$ in $d=2$ dimensions. To this end write the total density of states as a sum of a singular and a regular contribution as calculated in 14 b). (*Hint: Consider the derivative of the antiferromagnetic susceptibilities with respect to β and perform a Sommerfeld-like expansion for the regular part of the DOS.*)
- c) Discuss how the results of b) are modified in $d \geq 3$ dimensions.
- d) Consider the same electronic system for $d=2$, but now in presence of a local (repulsive) Hubbard interaction $U > 0$, and calculate within the random-phase approximation (RPA) the two (ferromagnetic and antiferromagnetic) susceptibilities in the corresponding particle-hole channel. On the basis of your RPA calculations, make your final considerations about the tendency of the system towards a given magnetic order at $T=0$.