12. The interaction with the Z-boson can be expressed in the form

$$\mathcal{L}_{NC} = \frac{g}{\cos \theta_W} J^{NC}_{\mu} Z^{\mu} \,, \tag{25}$$

where J^{NC}_{μ} is named the neutral current. The original relevant Lagrangian density of the GWS model is

$$\mathcal{L} = i\bar{L}\gamma_{\mu}D^{\mu}L + i\bar{R}\gamma_{\mu}D^{\mu}R\,, \qquad (26)$$

where the letters L and R stand for the left-handed dublet ψ_L and the right-handed singlet ψ_R of one fermion generation, e.g. (ν_e, e_L^-) and e_R^- , and the operator D_μ acts on L and R as

$$D^{\mu}L = \left(\partial^{\mu}\mathbb{1}_{2x2} - ig\,T^{a}A^{a\mu} - ig'\,\frac{Y}{2}\mathbb{1}_{2x2}B^{\mu}\right)L; \quad D^{\mu}R = \left(\partial^{\mu} - ig'\,\frac{Y}{2}B^{\mu}\right)R. \tag{27}$$

Extract from the Lagrangian density Eq. (26) the part

$$gJ^3_{\mu}A^{3\mu} + \frac{1}{2}g'J^Y_{\mu}B^{\mu}$$
(28)

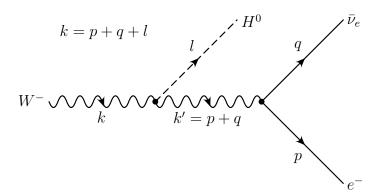
and perform the orthogonal transformation (22). Further use

$$J^{em}_{\mu} = J^3_{\mu} + \frac{1}{2} J^Y_{\mu} \tag{29}$$

and show that

$$J_{\mu}^{NC} = J_{\mu}^3 - \sin^2 \theta_W J_{\mu}^{em} \,. \tag{30}$$

13. Given is the decay of a real W boson, $W \to W^* + H^0 \to e \bar{\nu} + H^0$, where first a H^0 is radiated off.



The corresponding Lagrangians are

$$\mathcal{L}_{HWW} = gm_W H^0 W^+_\mu W^{\mu-} \tag{31}$$

$$\mathcal{L}^{CC} = \frac{g}{\sqrt{2}} J^{CC}_{\mu} W^{\mu +} + h.c.$$
 (32)

with $J_{\mu}^{CC} = \frac{1}{2} \bar{\nu}_e \gamma_{\mu} (1 - \gamma_5) e$. Note that the W^+ field operator annihilates a positively charged W boson or creates a negatively charged one. Calculate the matrix element

$$\mathcal{M} = \frac{ig^2 m_W}{2\sqrt{2}} \frac{1}{(p+q)^2 - m_W^2} \bar{u}(p)\gamma_\mu (1-\gamma_5)v(q) \,\epsilon^\mu \,. \tag{33}$$